# Thermogravimetric Study of the $\mathrm{Ln}_{2} \mathrm{O}_{3}-\mathrm{Co}-\mathrm{Co}_{2} \mathrm{O}_{3}$ System 

IV. $L n=$ La at 1100 and $1150^{\circ} \mathrm{C}$

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#### Abstract

Phase equilibria in the system $\mathrm{La}-\mathrm{Co}-\mathrm{O}$ at 1100 and $1150^{\circ} \mathrm{C}$ are established by changing the oxygen partial pressure from 0 to 12.00 in $-\log \left(P_{\mathrm{O}_{2}} / \mathrm{atm}\right)$, and a phase diagram of the $\mathrm{La}_{2} \mathrm{O}_{3}-\mathrm{Co}-\mathrm{Co}_{2} \mathrm{O}_{3}$ system is represented. Under the experimental conditions, $\mathrm{La}_{2} \mathrm{O}_{3}, \mathrm{CoO}, \mathrm{Co}, \mathrm{La}_{2} \mathrm{CoO}_{4}, \mathrm{LaCoO}_{3}$, and $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$ phases are present. Nonstoichiometry of the compound is presented with respect to the oxygen partial pressure, and activities of the components in the solid solutions are calculated by using the composition-oxygen partial pressure relations obtained. The standard Gibbs energy changes of the reactions in the phase diagrams are calculated and their linearity with reciprocal temperature is confirmed under the present experimental conditions. © 1997 Academic Press


The phase equilibrium in the $\mathrm{La}-\mathrm{Co}-\mathrm{O}$ system was established at $1130^{\circ} \mathrm{C}$ by Janecek and Wirtz (1), but their phase diagram was very rough, as illustrated in Fig. 1 in Ref. 1. Recently the phase equilibrium of the $\mathrm{La}-\mathrm{Co}-\mathrm{O}$ system at $1200^{\circ} \mathrm{C}$ was presented (2). In that report, $\mathrm{La}_{2} \mathrm{O}_{3}$, $\mathrm{Co}, \mathrm{CoO}, \mathrm{LaCoO}_{3}, \mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$, and $\mathrm{La}_{2} \mathrm{CoO}_{4}$ were stable and the standard Gibbs energy changes were obtained based on reactions in the phase diagram.

The objectives of the present study are (i) to establish detailed phase equilibria in the system at 1100 and $1150^{\circ} \mathrm{C}$ as a function of the oxygen partial pressures, (ii) to determine the thermochemical properties based on the phase equilibrium at different temperatures, and (iii) to ascertain whether the isothermal pattern of the phase diagram of the $\mathrm{La}-\mathrm{Co}-\mathrm{O}$ system changes with temperature.

## EXPERIMENTAL

Analytical grade $\mathrm{La}_{2} \mathrm{O}_{3}(99.9 \%)$ and $\mathrm{CoO}(99.9 \%)$ were used as starting materials. These oxides were dried by heating at $1100^{\circ} \mathrm{C}$ in air. Mixtures with desired ratios of $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}$ were prepared by thorough mixing in an agate mortar, calcined several times at $1100^{\circ} \mathrm{C}$ during the inter-
mediate mixing, and treated by the same procedures as those described previously (3).

Mixed gases of $\mathrm{CO}_{2}$ and $\mathrm{H}_{2}$ and of $\mathrm{CO}_{2}$ and $\mathrm{O}_{2}$ were used to obtain the oxygen partial pressures. The actual oxygen partial pressures of the gas phase were measured by means of a solid electrolytic cell of $\left(\mathrm{ZrO}_{2}\right)_{0.85}(\mathrm{CaO})_{0.15}(4)$.

The apparatus and procedures used to control the oxygen partial pressures and keep a constant temperature, the method of thermogravimetry, and the criterion for the establishment of equilibrium are the same as those described in a previous paper (5).

## RESULTS AND DISCUSSION

## (1) Phase Equilibria

(a) Co-O system. The $\mathrm{Co}-\mathrm{O}$ system is a subsystem of the present system. The system is reinvestigated at 1100 and $1150^{\circ} \mathrm{C}$ by using the present apparatus and procedures. It is known that in this system the compounds $\mathrm{CoO}, \mathrm{Co}_{3} \mathrm{O}_{4}$, and $\mathrm{Co}_{2} \mathrm{O}_{3}$ are present as stable phases, but under the present experimental conditions only the CoO phase is stable and it has a slight nonstoichiometric to oxygen-rich side. As shown in Table 1, for example, $x$ is $\sim 1.00$ and 1.011 in $\mathrm{CoO}_{x}$ at 10.33 and 0 in $-\log \left(P_{\mathrm{O}_{2}} / \mathrm{atm}\right)$ at $1100^{\circ} \mathrm{C}$. The weight in the atmosphere of $\log \left(P_{\mathrm{O}_{2}} /\right.$ atm $)=-12.00$ at which Co metal is stable is chosen as the standard weight.

The oxygen partial pressures in equilibrium with CoO and Co metal at 1100 and $1150^{\circ} \mathrm{C}$ are $10.33 \pm 0.03$ and $9.70 \pm 0.03$ in $-\log P_{\mathrm{O}_{2}}$, respectively.
(b) $\mathrm{La}_{2} \mathrm{O}_{3}-\mathrm{Co}-\mathrm{Co}_{2} \mathrm{O}_{3}$ system at $1100^{\circ} \mathrm{C}$. Nine samples with $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}$ mole ratios of $6 / 4,1,45 / 55,4 / 6,3 / 5,35 / 65$, $1 / 2,2 / 8$, and $1 / 9$ are prepared for use in thermogravimetry. In Fig. 1 the relationships between the oxygen partial pressure, $-\log \left(P_{\mathrm{O}_{2}} / \mathrm{atm}\right)$, on the ordinate and the weight change, $W_{\mathrm{O}_{2}} / W_{\mathrm{T}}$, on the abscissa are shown with four representative samples, 6/4 (Fig. 1a), 45/55 (Fig. 1b), 3/5 (Fig. 1c), and $2 / 8$ (Fig. 1d). Here, $W_{\mathrm{O}_{2}}$ is the weight increase

TABLE 1
Compositions, Symbols, Stability Ranges in Oxygen Partial Pressures, and Activities of Components in Solid Solutions

| Component | Compositions | Symbol | $\begin{gathered} -\log P_{\mathrm{O}_{2}} \\ (\mathrm{~atm}) \end{gathered}$ | $\log a_{i}$ |
| :---: | :---: | :---: | :---: | :---: |
| $1100^{\circ} \mathrm{C}$ |  |  |  |  |
| $\mathrm{La}_{2} \mathrm{CoO}_{4}$ | $\mathrm{La}_{2} \mathrm{CoO}_{4.00}$ | $A_{1}$ | 11.05 | 0 |
|  | $\mathrm{La}_{2} \mathrm{CoO}_{4.02}$ | $A_{2}$ | 10.33 | $\sim 0$ |
|  | $\mathrm{La}_{2} \mathrm{CoO}_{4.17}$ | $A_{3}$ | 5.70 | $-0.222$ |
|  | $\mathrm{La}_{2} \mathrm{CoO}_{4.20}$ | $A_{4}$ | 4.65 | -0.318 |
| $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$ | $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10.01}$ | $B_{1}$ | 5.70 | 0 |
|  | $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10.03}$ | $B_{2}$ | 4.95 | $-8 \times 10^{-3}$ |
|  | $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10.04}$ | $B_{3}$ | 4.65 | $-0.0137$ |
|  | $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10.10}$ | $B_{4}$ | 2.70 | -0.0848 |
| $\mathrm{LaCoO}_{3}$ | $\mathrm{LaCoO}_{2.92}$ | $C_{1}$ | 4.95 | 0 |
|  | $\mathrm{LaCoO}_{2.96}$ | $C_{2}$ | 2.70 | 0.0666 |
|  | $\mathrm{LaCoO}_{3.00}$ | $C_{3}$ | $0.680 \sim 0$ | 0.0919 |
| CoO | $\mathrm{CoO}_{1.00}$ | $D_{1}$ | 10.33 | 0 |
|  | $\mathrm{CoO}_{1.00}$ | $D_{2}$ | 5.70 | $\sim 0$ |
|  | $\mathrm{CoO}_{1.00}$ | $D_{3}$ | 4.95 | $\sim 0$ |
|  | $\mathrm{CoO}_{1.01}$ | $D_{4}$ | $1.50 \sim 0$ | $\sim 0$ |
| $1150^{\circ} \mathrm{C}$ |  |  |  |  |
| $\mathrm{La}_{2} \mathrm{CoO}_{4}$ | $\mathrm{La}_{2} \mathrm{CoO}_{4.01}$ | $A_{1}^{\prime}$ | 10.50 | 0 |
|  | $\mathrm{La}_{2} \mathrm{CoO}_{4.02}$ | $A_{2}^{\prime}$ | 9.70 | $-7.44 \times 10^{-3}$ |
|  | $\mathrm{La}_{2} \mathrm{CoO}_{4.17}$ | $A_{3}^{\prime}$ | 5.00 | -0.189 |
|  | $\mathrm{La}_{2} \mathrm{CoO}_{4.20}$ | $A_{4}^{\prime}$ | 3.80 | $-0.212$ |
| $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$ | $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{9.93}$ | $B_{1}^{\prime}$ | 5.00 | 0 |
|  | $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{9.94}$ | $B_{2}^{\prime}$ | 4.30 | 0.0218 |
|  | $\mathrm{La}_{4} \mathrm{CoO}_{9.96}$ | $B_{3}^{\prime}$ | 3.80 | -0.0381 |
|  | $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{9.99}$ | $B_{4}^{\prime}$ | 2.20 | 0.0561 |
| $\mathrm{LaCoO}_{3}$ | $\mathrm{LaCoO}_{2.91}$ | $C_{1}^{\prime}$ | 4.50 | 0 |
|  | $\mathrm{LaCoO}_{2.97}$ | $C_{2}^{\prime}$ | 2.20 | 0.0726 |
|  | $\mathrm{LaCoO}_{3.00}$ | $C_{3}^{\prime}$ | $0.68 \sim 0$ | 0.0819 |
| CoO | $\mathrm{CoO}_{1.00}$ | $D_{1}^{\prime}$ | 9.70 | 0 |
|  | $\mathrm{CoO}_{1.00}$ | $D_{2}^{\prime}$ | 5.00 | $\sim 0$ |
|  | $\mathrm{CoO}_{1.00}$ | $D_{3}^{\prime}$ | 4.30 | $\sim 0$ |
|  | $\mathrm{CoO}_{1.01}$ | $D_{4}^{\prime}$ | $4.30 \sim 0$ | $\sim 0$ |

of the samples from the reference weight at $\log P_{\mathrm{O}_{2}}=$ -12.00 , at which $\mathrm{La}_{2} \mathrm{O}_{3}$ and Co are stable, and $W_{\mathrm{T}}$ is the total weight gain from the reference state to the weight at $1 \mathrm{~atm} \mathrm{O}_{2}$, at which $\mathrm{La}_{2} \mathrm{O}_{3}+\mathrm{LaCoO}_{3}$ or $\mathrm{LaCoO}_{3}+\mathrm{CoO}$ are stable depending on the total composition of the samples. Abrupt weight changes are observed at 11.05, 10.33, $5.70,4.95,4.65$, and 2.70 in $-\log P_{\mathrm{O}_{2}}$ in Fig. 1. These values correspond to the oxygen partial pressure in equilibrium with three solid phases.
The compositions of the ternary compounds, symbols, stability ranges in $\log P_{\mathrm{O}_{2}}$, and activities of the components in solid solutions are shown in Table 1.

Based on the above thermogravimetric results, a phase diagram was drawn and is illustrated in Fig. 2 as $\mathrm{La}_{2} \mathrm{O}_{3}-\mathrm{Co}-\mathrm{Co}_{2} \mathrm{O}_{3}$ although $\mathrm{Co}_{2} \mathrm{O}_{3}$ is not stable under the experimental conditions. The numerical values in the three solid fields in Fig. 2 are the six values in $-\log P_{\mathrm{O}_{2}}$ described above. Six phases, $\mathrm{La}_{2} \mathrm{O}_{3}, \mathrm{Co}, \mathrm{La}_{2} \mathrm{CoO}_{4}, \mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$,
$\mathrm{LaCoO}_{3}$, and CoO , are stable under the present conditions. The stable phases are the same as those at $1200^{\circ} \mathrm{C}$ (2). $\mathrm{Co}_{3} \mathrm{O}_{4}$ is not found. Except for Co and $\mathrm{La}_{2} \mathrm{O}_{3}$, the other phases have nonstoichiometric compositions. The relationship between the composition of a solid solution and the oxygen partial pressure is represented by a linear equation such as $N_{\mathrm{O}} / N_{\text {component }}=a \log P_{\mathrm{O}_{2}}+b$, which can be used to calculate the activity of the components in the solid solutions with the Gibbs-Duhem equation (6). Here, $N_{o}$ and $N_{\text {component }}$ are the mole fraction of oxygen and component (cf. Table 1) in the solid solutions, respectively. The $a$ and $b$ values for each solid solution, which are obtained from the thermogravimetric results, are tabulated in Table 2.

As for the oxygen deficiency of the $\mathrm{LaCoO}_{3}$ phase, Seppanen et al. (7) expressed the deficiency by a general equation, $\mathrm{LaCoO}_{3-d}$, as $d=k P_{\mathrm{O}_{2}}^{-1 / 2}$ and at 1300 K reaches the value 0.091 at the phase boundary $\mathrm{LaCoO}_{3-d^{-}}$ $\mathrm{CoO}-\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$. The present value $0.08\left(\mathrm{LaCoO}_{2.92}\right.$ $\left(C_{1}\right)$ ) in Table 1 might be reasonable considering the difference in temperature.

To express the chemical composition of the three ternary compounds, Janecek et al. (1) presented the general formula $\mathrm{La}_{n+1} \mathrm{Co}_{n} \mathrm{O}_{3 n+1}\left(\mathrm{LaO}\left(\mathrm{LaCoO}_{3}\right)_{n}\right)$ as the perovskite-related layered structure. The phase diagram suggests that the compositions $\mathrm{LaCoO}_{3}, \mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$, and $\mathrm{La}_{2} \mathrm{CoO}_{4}$ can be represented by the compositions $\mathrm{La}_{2} \mathrm{O}_{3} \cdot \mathrm{Co}_{2} \mathrm{O}_{3}$, $2 \mathrm{La}_{2} \mathrm{O}_{3} \cdot \mathrm{Co}_{3} \mathrm{O}_{4}$, and $\mathrm{La}_{2} \mathrm{O}_{3} \cdot \mathrm{CoO}$, respectively, that is, a two-component system of $\mathrm{La}_{2} \mathrm{O}_{3}$-cobalt oxide, although there is no crystallographic evidence for it now.

Recently, Lewandowski et al. (8) reported that the composition of $\mathrm{La}_{2} \mathrm{CoO}_{4}$ did not exist in the range of temperatures and $P_{\mathrm{O}_{2}}$ investigated and proposed a lanthanum cation-deficient composition, $\mathrm{La}_{1.83} \mathrm{CoO}_{4}$. As shown in Fig. 2 in the present case, lanthanum-cation deficiency is not found. Both lines, which originate from Co and $\mathrm{La}_{2} \mathrm{O}_{3}$ apices and go through the experimental points, join at $A_{1}$ and have the composition $20 \mathrm{~mol} \% \mathrm{Co}, 60 \mathrm{~mol} \%_{\mathrm{La}_{2} \mathrm{O}_{3}}$, and $20 \mathrm{~mol} \% \mathrm{Co}_{2} \mathrm{O}_{3}$, that is, $\mathrm{La}_{2} \mathrm{CoO}_{4}$. The result is the same as those at $1200^{\circ} \mathrm{C}$.
(c) $\mathrm{La}_{2} \mathrm{O}_{3}-\mathrm{Co}-\mathrm{Co}_{2} \mathrm{O}_{3}$ system at $1150^{\circ} \mathrm{C}$. Five samples with $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}$ mole ratios of $6 / 4,1,4 / 6,1 / 2$, and $2 / 8$ were prepared for use in thermogravimetry, although they were anticipated to have the same pattern as those at 1100 and $1200^{\circ} \mathrm{C}$. In Fig. 1 the relationships between the oxygen partial pressure and weight changes are shown for four representative samples, $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}$ mole ratios of $6 / 4$ (Fig. 1e), 1 (Fig. 1f), $4 / 6$ (Fig. 1g), and $2 / 8$ (Fig. 1h). The reference weight at $\log P_{\mathrm{O}_{2}}=-12.00$ was chosen here also. Based on Figs. 1e-1h, a pattern the same as that in Fig. 2 at $1100^{\circ} \mathrm{C}$ can be easily seen. The oxygen partial pressures in equilibrium with three solid phases and the activities of the component in the solid solutions are also given in Table 1.


FIG. 1. The relationship between the oxygen partial pressure, $\log \left(P_{\mathrm{O}_{2}} /\right.$ atm $)$, and the weight change of the samples, $W_{\mathrm{O}_{2}} / W_{\mathrm{T}}$. At $1100^{\circ} \mathrm{C}$ : (a) $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}=6 / 4$, (b) $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}=45 / 55$, (c) $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}=3 / 5$, and (d) $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}=2 / 8$. At $1150^{\circ} \mathrm{C}:(\mathrm{e}) \mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}=6 / 4$, (f) $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}=5 / 5$, (g) $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}=4 / 6$, and (h) $\mathrm{La}_{2} \mathrm{O}_{3} / \mathrm{CoO}=2 / 8$.


FIG. 1-Continued


FIG. 2. Phase equilibrium in the $\mathrm{La}_{2} \mathrm{O}_{3}-\mathrm{Co}-\mathrm{Co}_{2} \mathrm{O}_{3}$ system at $1100^{\circ} \mathrm{C}$. Numerical values in the three phase regions are the oxygen partial pressures in $-\log \left(P_{\mathrm{O}_{2}} / \mathrm{atm}\right)$ in equilibrium with three solid phases. Abbreviations are the same as those in Table 1.

## (2) The Standard Gibbs Energy Change of Reaction and Its Dependency on Temperature

On the basis of the established phase diagram, the standard Gibbs energy changes of reactions, which appear in the

TABLE 2 $\boldsymbol{a}$ and $\boldsymbol{b}$ Values of Solid Solution

| Component | $a$ | $b$ |
| :--- | :---: | :---: |
|  | $1100^{\circ} \mathrm{C}$ |  |
| $\mathrm{La}_{2} \mathrm{CoO}_{4}$ | 0.0313 | 0.345 |
| $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$ | 0.0313 | 0.188 |
| $\mathrm{LaCoO}_{3}$ | 0.0160 | 0.002 |
|  |  |  |
| $\mathrm{La}_{2} \mathrm{CoO}_{4}$ | $1150^{\circ} \mathrm{C}$ |  |
| $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$ | 0.0151 | 0.171 |
| $\mathrm{LaCoO}_{3}$ | 0.0210 | 0.0355 |

$$
\text { Note. } N_{\mathrm{O}} / N_{\text {component }}=a \log P_{\mathrm{O}_{2}}+b
$$

diagram and are shown in Table 3, are determined with the equation $\Delta G^{\circ}=-R T \ln K$. Here, $R$ is the gas constant, $T$ the absolute temperature, and $K$ the equilibrium constant of the reaction. Activities of the components in solid solutions, which are necessary to calculate $\Delta G^{\circ}$ values, are given in Table 1. The standard state of the components in solid solution can be chosen arbitrarily in each solid solution as $\log a_{i}=0$ in Table 1.
$\Delta G^{\circ}$ values obtained for each reaction are given in Table 3 together with those at $1200^{\circ} \mathrm{C}$ (1). Assuming that the activity of the components in the solid solutions is unity, $\Delta G^{\circ}$ values for reactions [3]-[5] in Table 3 are -17.7, -74.9 , and $-61.1 \mathrm{~kJ} \mathrm{~mol}^{-1}$, at $1100^{\circ} \mathrm{C}$, and -15.0 , -70.5 , and $-51.8 \mathrm{~kJ} \mathrm{~mol}^{-1}$ at $1150^{\circ} \mathrm{C}$. The differences in

TABLE 3
The Standard Gibbs Energy Change of Reactions

| Reaction | Temperature $\left({ }^{\circ} \mathrm{C}\right)$ | $\begin{gathered} -\log P_{\mathrm{O}_{2}} \\ (\operatorname{atm}) \end{gathered}$ | $\begin{gathered} -\Delta G^{\circ} \\ (\mathrm{kJ}) \end{gathered}$ |
| :---: | :---: | :---: | :---: |
| [1] $\mathrm{Co}+\mathrm{La}_{2} \mathrm{O}_{3}+\frac{1}{2} \mathrm{O}_{2}=\mathrm{La}_{2} \mathrm{CoO}_{4}$ | 1100 | 11.05 | 145.3 |
|  | 1150 | 10.50 | 143.1 |
|  | 1200 | 10.02 | 141.1 |
| [2] $\mathrm{Co}+\frac{1}{2} \mathrm{O}_{2}=\mathrm{CoO}$ | 1100 | 10.33 | 135.8 |
|  | 1150 | 9.70 | 132.2 |
|  | 1200 | 9.30 | 131.1 |
| [3] $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}+\frac{1}{4} \mathrm{O}_{2}=\frac{1}{2} \mathrm{La}_{2} \mathrm{O}_{3}+3 \mathrm{LaCoO}_{3}$ | 1100 | 2.70 | 25.2 |
|  | 1150 | 2.20 | 19.4 |
|  | 1200 | 1.87 | 17.5 |
| [4] $2 \mathrm{La}_{2} \mathrm{CoO}_{4}+\mathrm{CoO}+\frac{1}{2} \mathrm{O}_{2}=\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$ | 1100 | 5.70 | 86.6 |
|  | 1150 | 5.00 | 76.2 |
|  | 1200 | 4.70 | 70.1 |
| [5] $3 \mathrm{La}_{2} \mathrm{CoO}_{4}+\frac{1}{2} \mathrm{O}_{2}=\mathrm{La}_{2} \mathrm{O}_{3}+\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}$ | 1100 | 4.65 | 85.8 |
|  | 1150 | 3.80 | 70.1 |
|  | 1200 | 3.05 | 57.8 |
| [6] $\mathrm{La}_{4} \mathrm{Co}_{3} \mathrm{O}_{10}+\mathrm{CoO}+\frac{1}{2} \mathrm{O}_{2}=4 \mathrm{LaCoO}_{3}$ | 1100 | 4.95 | 65.1 |
|  | 1150 | 4.30 | 58.6 |
|  | 1200 | 3.90 | 52.3 |

Note. Values at $1200^{\circ} \mathrm{C}$ are quoted from Ref. 2.


Fig. 3. The relationship between Gibbs energy change $\left(-\Delta G^{\circ}\right)$ and reciprocal temperature $\left(1 / T \times 10^{4}\right)$. Numbers, $1-6$ in the figures correspond to the reactions in Table 3.
these values are significant. As the phase equilibria are not known exactly, Seppanen et al. (7) did not use activity values in the calculation of Gibbs energy changes.

The combination of reactions given in Table 3 can form new equations, for example, reaction [7], $\mathrm{La}_{2} \mathrm{O}_{3}+\mathrm{CoO}$ $=\mathrm{La}_{2} \mathrm{CoO}_{4}([1]-[2])$; and [8] $\mathrm{CoO}+1 / 2 \mathrm{La}_{2} \mathrm{O}_{3}+1 / 4 \mathrm{O}_{2}$ $=\mathrm{LaCoO}_{3}([6]-[3]) . \Delta G^{\circ}$ values are easily obtained as -9.5 and $-10.9 \mathrm{~kJ} \mathrm{~mol}^{-1}$ for [7] at 1100 and $1150^{\circ} \mathrm{C}$ and -40.1 and $-38.6 \mathrm{~kJ} \mathrm{~mol}^{-1}$ for [8] at 1100 and $1150^{\circ} \mathrm{C}$, respectively. For reaction [7], Sreedharan et al. [9] presented the equation $\Delta G^{\circ}\left(\mathrm{kJ} \mathrm{mol}{ }^{-1}\right)=12.24-0.0149 T$ ( $973-1373 \mathrm{~K}, \pm 1.02 \mathrm{~kJ}$ ). The $\Delta G^{\circ}$ value for reaction [7] at $1100^{\circ} \mathrm{C}$ is -8.2 kJ and by extrapolating the equation to $1150^{\circ} \mathrm{C}$ we obtain a $\Delta G^{\circ}$ value of -9.0 kJ . Both values are in fairly good agreement with the present results.

As for reaction $[2] \times 2, \Delta G^{\circ}\left(\mathrm{cal} \mathrm{mol}^{-1}\right)=-111,800$ $+33.8 T( \pm 2 \mathrm{kcal}, 298-1400 \mathrm{~K})$ was reported (10). Values of -136.9 kJ at $1100^{\circ} \mathrm{C}$ and -129.8 kJ at $1150^{\circ} \mathrm{C}$ are obtained from the equation. These values are in good agreement with the present -135.8 and -132.2 kJ , respectively.

In Fig. 3 the standard Gibbs energy change versus reciprocal temperature is plotted. A linear relationship between both variables is observed.

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